

Painlevé equations, Cherednik algebras and q -Askey scheme

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Painlevé equations are **non linear** second order ODE

$$\frac{d^2 w}{dz^2} = F\left(z, w, \frac{dw}{dz}\right), \quad z \in \overline{\mathbb{C}},$$

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- 1 **Painlevé–Kowalevski property:** $w(z; c_1, c_2)$ have no *critical points* that depend on c_1, c_2 .
- 2 For generic c_1, c_2 , $w(z; c_1, c_2)$ are **new** functions, **Painlevé Transcendents**.

$$\frac{d^2 w}{dz^2} = 6w^2 + z, \quad \frac{d^2 w}{dz^2} = 2w^3 + zw + \alpha$$

$$\frac{d^2 w}{dz^2} = \frac{1}{w} \left(\frac{dw}{dz} \right)^2 - \frac{1}{z} \frac{dw}{dz} + \frac{\alpha w^2 + \beta}{z} + \gamma w^3 + \frac{\delta}{w}$$

$$\frac{d^2 w}{dz^2} = \frac{1}{2w} \left(\frac{dw}{dz} \right)^2 + \frac{3}{2} w^3 + 4zw^2 + 2(z^2 - \alpha)w + \frac{\beta}{w}$$

$$\begin{aligned} \frac{d^2 w}{dz^2} = & \left(\frac{1}{2w} + \frac{1}{w-1} \right) \left(\frac{dw}{dz} \right)^2 - \frac{1}{z} \frac{dw}{dz} + \frac{\gamma w}{z} + \\ & + \frac{(w-1)^2}{z^2} \left(\alpha w + \frac{\beta}{w} \right) + \frac{\delta w(w+1)}{w-1} \end{aligned}$$

$$\begin{aligned} \frac{d^2 w}{dz^2} = & \frac{1}{2} \left(\frac{1}{w} + \frac{1}{w-1} + \frac{1}{w-z} \right) w_z^2 - \left(\frac{1}{z} + \frac{1}{z-1} + \frac{1}{w-z} \right) w_z + \\ & + \frac{w(w-1)(w-z)}{z^2(z-1)^2} \left[\alpha + \beta \frac{z}{w^2} + \gamma \frac{z-1}{(w-1)^2} + \delta \frac{z(z-1)}{(w-z)^2} \right] \end{aligned}$$

Painlevé transcendents in Random Matrix Theory:

Example: (S_N, P_N) group of permutations π of $1, 2, \dots, N$ with uniform probability distribution $P_N(\pi) = \frac{1}{N!}$.

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$$\lim_{N \rightarrow \infty} \text{Prob}\left(\frac{l_N(\pi) - 2\sqrt{N}}{N^{1/6}} \leq t\right) = \exp\left\{-\int_t^\infty (t-z)w^2(z) \, dz\right\}$$

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Only special solutions appear in this way

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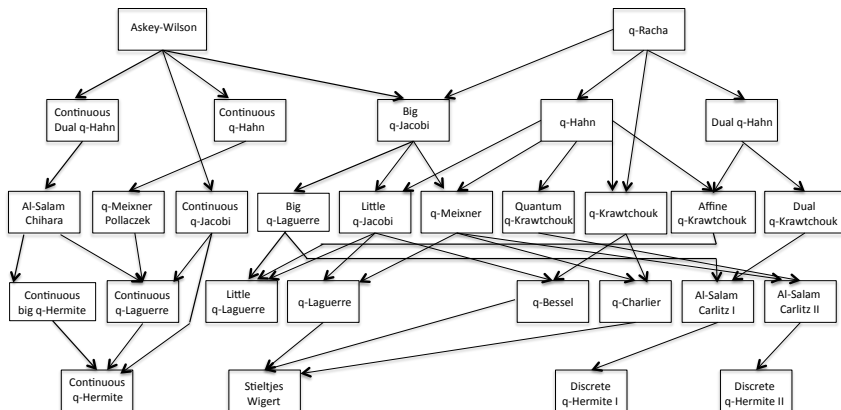
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- To each Painlevé equation we associate a cubic surface called *monodromy manifold*.
- The monodromy manifold admits a Poisson bracket and can be quantised.
- These quantum algebras admit representations on the space of Laurent polynomials.
- The eigenspaces in this representation are spanned by q -orthogonal polynomials (q -Askey scheme).

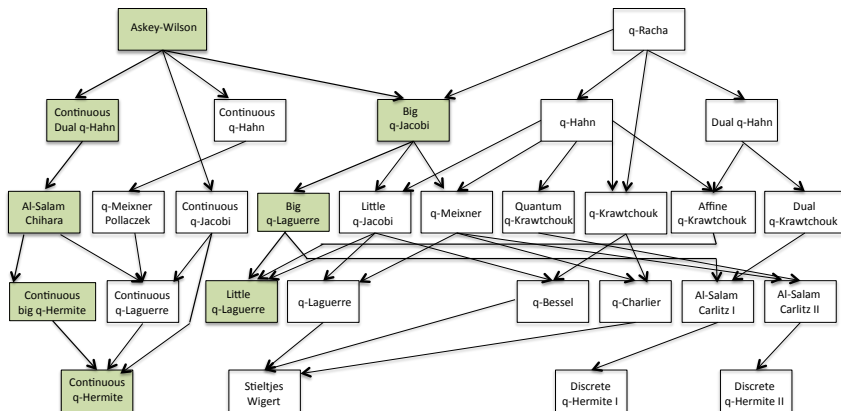
q-Askey scheme

Koekoek, Lesky, Swarttouw 2010

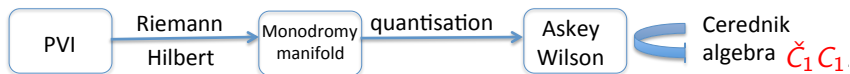


Painlevé equations and q -Askey polynomials

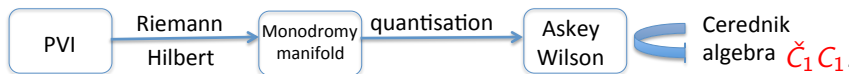
M.M. arXiv:1307.6140



Other results

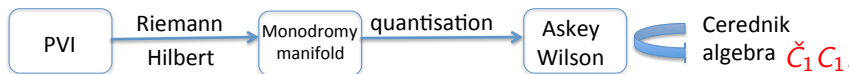


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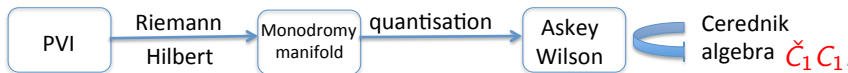
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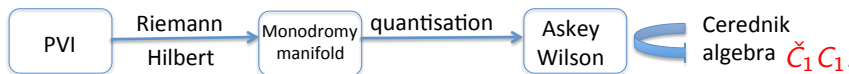
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- Confluence Painlevé equations,
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- Confluence Askey Wilson algebra,
- Confluence Cherednik algebra $\Rightarrow$ **seven new algebras.** M.M.

arXiv:1307.6140

Askey Wilson polynomials

q -Pochhammer symbol: $(a; q)_k := \prod_{j=0}^{k-1} (1 - a q^j)$

$(a_1, \dots, a_r; q)_k := (a_1; q)_k (a_2; q)_k \cdots (a_r; q)_k$

q -hypergeometric function:

$${}_r\varphi_{r-1} \left(\begin{matrix} a_1, a_2, \dots, a_r \\ b_1, \dots, b_{r-1} \end{matrix} ; q, x \right) := \sum_{k=0}^{\infty} \frac{(a_1, a_2, \dots, a_r; q)_k}{(b_1, \dots, b_{r-1}; q)_k} x^k.$$

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Askey-Wilson polynomials:

$$\begin{aligned} p_n(z; a, b, c, d \mid q) &:= {}_4\varphi_3 \left(\begin{matrix} q^{-n}, q^{n-1}abcd, az, az^{-1} \\ ab, ac, ad \end{matrix} ; q, q \right) \\ &= \sum_{k=0}^{n+1} \frac{(q^{-n}, q^{n-1}abcd, az, az^{-1}; q)_k}{(ab, ac, ad; q)_k} q^k. \end{aligned}$$

Facts about Askey Wilson polynomials

- They are special cases of Macdonald polynomials.
- Their algebraic structure is controlled by the Cherednik algebra of type $\check{C}_1 C_1$.
- They contain all other basic hypergeometric polynomials as degenerations.

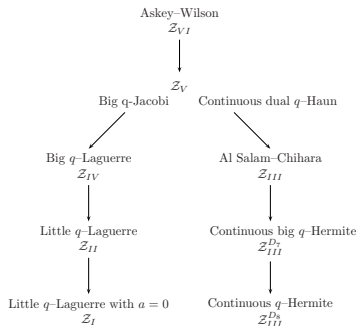
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Open question: what algebras control all other elements of the q -Askey scheme?

Quantum algebra results

- For each Painlevé equation, a new quantum algebra and its representation theory is introduced.
- These algebras control some elements of the q -Askey scheme.



Cherednik algebra of type $\check{C}_1 C_1$

Cherednik '92, Sahi '99

Algebra generated by $V_0, V_1, \check{V}_0, \check{V}_1$:

$$(V_0 - k_0)(V_0 + k_0^{-1}) = 0$$

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$$(\check{V}_0 - u_0)(\check{V}_0 + u_0^{-1}) = 0$$

$$(\check{V}_1 - u_1)(\check{V}_1 + u_1^{-1}) = 0$$

$$\check{V}_1 V_1 V_0 \check{V}_0 = q^{-1/2},$$

k_0, k_1, u_0, u_1 scalars.

Representation theory:

Cherednik algebra on Laurent polynomials

$$V_1[f] = \frac{z \left[u_1 z^2 - \frac{1}{u_1} + \left(k_1 - \frac{1}{k_1} \right) z \right]}{z^2 - 1} f\left(\frac{1}{z}\right) - \frac{z \left(u_1 - \frac{1}{u_1} \right) + k_1 - \frac{1}{k_1}}{z^2 - 1} f(z)$$

$$V_0[f] = \frac{q k_0 - \frac{z^2}{k_0} + \sqrt{q} \left(u_0 - \frac{1}{u_0} \right) z}{z^2 - 1} f\left(\frac{q}{z}\right) - \frac{z \left[\left(k_0 - \frac{1}{k_0} \right) + \sqrt{q} \left(u_0 - \frac{1}{u_0} \right) \right]}{z^2 - 1} f(z)$$

The operator $\frac{1 + \check{V}_1}{1 + u_1^2}$ acts as symmetrizer $\Rightarrow$ representation of the spherical sub-algebra on **symmetric** Laurent polynomials.

Spherical sub-algebra:

generated by: $e = \frac{1+\check{V}_1}{1+u_1^2},$

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$$X_2 = e (\check{V}_1 V_0 + (\check{V}_1 V_0)^{-1}),$$

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$$q^{-1/2} X_1 X_2 - q^{1/2} X_2 X_1 = (q^{-1} - q) X_3 + (q^{-1/2} - q^{1/2}) \omega_3 e$$

$$q^{-1/2} X_2 X_3 - q^{1/2} X_3 X_2 = (q^{-1} - q) X_1 + (q^{-1/2} - q^{1/2}) \omega_1 e$$

$$q^{-1/2} X_3 X_1 - q^{1/2} X_1 X_3 = (q^{-1} - q) X_2 + (q^{-1/2} - q^{1/2}) \omega_2 e$$

$$q^{\frac{1}{2}} X_2 X_1 X_3 - q X_2^2 - q^{-1} X_1^2 - q X_3^2 + q^{\frac{1}{2}} \omega_2 X_2 + q^{-\frac{1}{2}} \omega_1 X_1 + q^{\frac{1}{2}} \omega_3 X_3 = \omega_4 e.$$

Zhedanov '91, Oblomkov '04

Spherical sub-algebra and Askey Wilson polynomials

$$X_1(p_n(z, a, b, c, d|q)) = (z + \frac{1}{z})p_n(z, a, b, c, d|q),$$

$$X_2(p_n(z, a, b, c, d|q)) = (q^{-n} + abcdq^{n-1})(p_n(z, a, b, c, d|q)).$$

Koornwinder '07

Recap:

- Cherednik algebra of type $\check{C}_1 C_1$:

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Represented on
Laurent polynomials
Non-symmetric
Askey-Wilson

- Spherical subalgebra (Zhedanov):

$$q^{-1/2} X_1 X_2 - q^{1/2} X_2 X_1 = (q^{-1} - q) X_3 + (q^{-1/2} - q^{1/2}) \omega_3 e,$$

$$q^{-1/2} X_2 X_3 - q^{1/2} X_1 X_3 = (q^{-1} - q) X_1 + (q^{-1/2} - q^{1/2}) \omega_1 e,$$

$$q^{-1/2} X_3 X_1 - q^{1/2} X_1 X_3 = (q^{-1} - q) X_2 + (q^{-1/2} - q^{1/2}) \omega_2 e.$$

Representation spanned by Askey-Wilson polynomials

All Painlevé equations are **isomonodromic deformation equations** (Jimbo-Miwa 1980)

$$\frac{dB}{d\lambda} - \frac{dA}{dz} = [A, B]$$

$$A = A(\lambda; z, w, w_z), \quad B = B(\lambda; z, w, w_z) \in \mathfrak{sl}_2.$$

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This means that **the monodromy data** of the linear system

$$\frac{dY}{d\lambda} = A(\lambda; z, w, w_z)Y$$

are **locally constant along solutions of the Painlevé equation**.

The monodromy data play the role of initial conditions

PVI monodromy data

$$M_1, M_2, M_3, M_\infty \in SL(2, \mathbb{C}) \text{ s.t.}$$

$$\text{eigenv}(M_i) = \exp(\pm i p_i), \quad i = 1, 2, 3, \infty,$$

$$M_\infty M_1 M_2 M_3 = 1.$$

Jimbo, Miwa '81

Cherednik algebra as quantisation of the monodromy group

$$(M_3 - e^{ip_3})(M_3 - e^{-ip_3}) = 0,$$

$$(M_2 - e^{ip_2})(M_2 - e^{-ip_2}) = 0,$$

$$(M_1 - e^{ip_1})(M_1 - e^{-ip_1}) = 0,$$

$$(M_\infty - e^{ip_\infty})(M_\infty - e^{-ip_\infty}) = 0$$

$$M_\infty M_3 M_2 M_1 = 1.$$

Cherednik algebra as quantisation of the monodromy group

$$\begin{aligned}
 (M_3 - e^{ip_3})(M_3 - e^{-ip_3}) &= 0, & (V_0 - k_0)(V_0 + k_0^{-1}) &= 0, \\
 (M_2 - e^{ip_2})(M_2 - e^{-ip_2}) &= 0, & (V_1 - k_1)(V_1 + k_1^{-1}) &= 0, \\
 (M_1 - e^{ip_1})(M_1 - e^{-ip_1}) &= 0, & (\check{V}_1 - u_1)(\check{V}_1 + u_1^{-1}) &= 0, \\
 (M_\infty - e^{ip_\infty})(M_\infty - e^{-ip_\infty}) &= 0 & (\check{V}_0 - u_0)(\check{V}_0 + u_0^{-1}) &= 0, \\
 M_\infty M_3 M_2 M_1 &= 1. & \check{V}_1 V_1 V_0 \check{V}_0 &= q^{-1/2}.
 \end{aligned}$$

$$M_1 \rightarrow i\check{V}_1, \quad M_2 \rightarrow iV_1, \quad M_3 \rightarrow iV_0, \quad M_\infty \rightarrow i\check{V}_0.$$

$$u_1 = -ie^{-\frac{p_1}{2}}, \quad k_0 = -ie^{-\frac{p_3}{2}}, \quad k_1 = -ie^{-\frac{p_2}{2}}, \quad u_0 = -ie^{-S_1 - S_2 - S_3}.$$

Monodromy manifold:

$$x_1 = \text{Tr}(M_2 M_3), \quad x_2 = \text{Tr}(M_2 M_3), \quad x_3 = \text{Tr}(M_3 M_1),$$

$$x_1 x_2 x_3 + x_1^2 + x_2^2 + x_3^2 + \omega_1 x_1 + \omega_2 x_2 + \omega_3 x_3 = \omega_4$$

Poisson bracket:

$$\{x_1, x_2\} = x_1 x_2 + 2x_3 + \omega_3, \quad \{x_2, x_3\} = x_2 x_3 + 2x_1 + \omega_1,$$

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Monodromy manifolds for the Painlevé equations

$$PVI \quad x_1 x_2 x_3 + x_1^2 + x_2^2 + x_3^2 + \omega_1 x_1 + \omega_2 x_2 + \omega_3 x_3 = \omega_4$$

$$PV \quad x_1 x_2 x_3 + x_1^2 + x_2^2 + \omega_1 x_1 + \omega_2 x_2 + \omega_3 x_3 = \omega_4$$

$$PIV \quad x_1 x_2 x_3 + x_1^2 + \omega_1 x_1 + \omega_2 x_2 + \omega_2 x_3 + 1 = \omega_4$$

$$PIII \quad x_1 x_2 x_3 + x_1^2 + x_2^2 + \omega_1 x_1 + \omega_2 x_2 = \omega_1 - 1$$

$$PII \quad x_1 x_2 x_3 + x_1 + x_2 + x_3 = \omega_4$$

$$PI \quad x_1 x_2 x_3 + x_1 + x_2 + 1 = 0$$

Poisson structure on the monodromy manifolds

Cubic surface $M_\varphi := \mathbb{C}[x_1, x_2, x_3] / \langle \varphi(x_1, x_2, x_3) = 0 \rangle$

Poisson bracket:

$$\{x_1, x_2\} = \frac{\partial \varphi}{\partial x_3}, \quad \{x_2, x_3\} = \frac{\partial \varphi}{\partial x_1}, \quad \{x_3, x_1\} = \frac{\partial \varphi}{\partial x_2}.$$

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Example PV: $\varphi := x_1 x_2 x_3 + x_1^2 + x_2^2 + \omega_1 x_1 + \omega_2 x_2 + \omega_3 x_3 - \omega_4,$

$$\begin{aligned} \{x_1, x_2\} &= x_1 x_2 + \omega_3, & \{x_2, x_3\} &= x_2 x_3 + 2x_1 + \omega_1, \\ \{x_3, x_1\} &= x_1 x_3 + 2x_2 + \omega_2. \end{aligned}$$

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Quantisation (Teichmüller theory and cluster algebra)

(L. Chekhov and M.M. J.Phys A 2010).

For $d \rightarrow 0$ ($k_0 \rightarrow 0$, $u_0 \rightarrow \epsilon u_0$):

- First confluent Cherednik algebra:

$$\begin{aligned}
 V_0^2 + V_0 &= 0, \\
 (V_1 - k_1)(V_1 + k_1^{-1}) &= 0 \\
 \check{V}_0^2 + u_0^{-1} \check{V}_0 &= 0, \\
 (\check{V}_1 - u_1)(\check{V}_1 + u_1^{-1}) &= 0 \\
 q^{1/2} \check{V}_1 V_1 V_0 &= \check{V}_0 + u_0^{-1}, \\
 q^{1/2} \check{V}_0 \check{V}_1 V_1 &= V_0 + 1.
 \end{aligned}$$

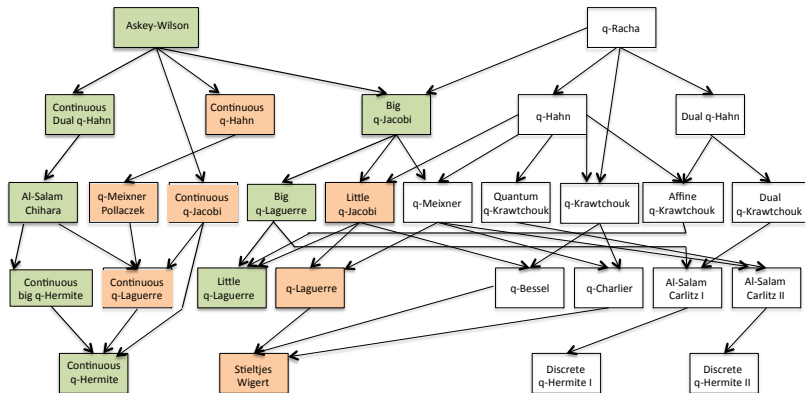
Represented on
Laurent polynomials
Non-symmetric
Continuous
dual q-Hahn
polynomials (M.M. arXiv:1409.4287)

- Spherical subalgebra (confluent Zhedanov):

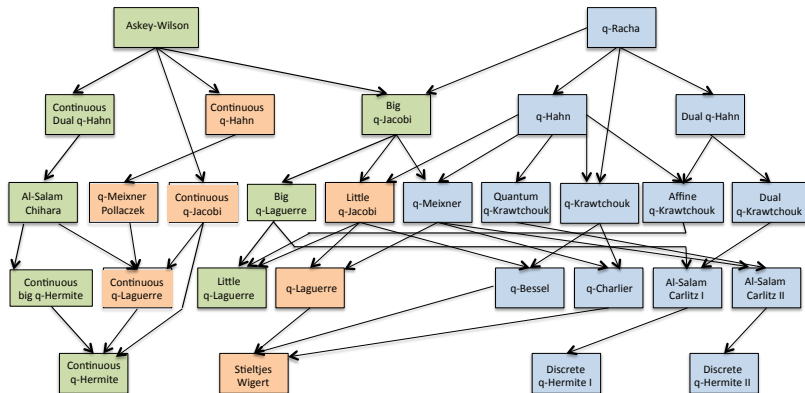
$$\begin{aligned}
 q^{-1/2} X_1 X_2 - q^{1/2} X_2 X_1 &= (q^{-1/2} - q^{1/2}) \omega_3 e, \\
 q^{-1/2} X_2 X_3 - q^{1/2} X_1 X_3 &= (q^{-1} - q) X_1 + (q^{-1/2} - q^{1/2}) \omega_1 e, \\
 q^{-1/2} X_3 X_1 - q^{1/2} X_1 X_3 &= (q^{-1} - q) X_2 + (q^{-1/2} - q^{1/2}) \omega_2 e.
 \end{aligned}$$

Representation spanned by Continuous dual q-Hahn
polynomials

Outlook



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- q -difference and elliptic Painlevé equations may correspond to elliptic hypergeometric bi-orthogonal polynomials.
- Multivariable high order analogues of the Painlevé equations $\rightarrow$ *confluence scheme for Macdonald polynomials*